NOTE ON 2D SCHRÖDINGER OPERATORS WITH δ -INTERACTIONS ON ANGLES AND CROSSING LINES

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In this note we sharpen the lower bound previously obtained by Lobanov et al [LLP10] for the spectrum of the 2D Schrödinger operator with a δ -interaction supported on a planar angle. Using the same method we obtain the lower bound on the spectrum of the 2D Schrödinger operator with a δ -interaction supported on crossing straight lines. The latter operators arise in the three-body quantum problem with δ -interactions between particles.

Keywords: Schrödinger operators, δ -interactions, spectral estimates.

1. Introduction

Self-adjoint Schrödinger operators with δ -interactions supported on sufficiently regular hypersurfaces can be defined via closed, densely defined, symmetric and lower-semibounded quadratic forms using the first representation theorem, see [BEKS94], [BLL13].

 δ -interactions on angles. In our first model the support of the δ -interaction is the set $\Sigma_{\varphi} \subset \mathbb{R}^2$, which consists of two rays meeting at the common origin and constituting the angle $\varphi \in (0, \pi]$ as in Figure 1.

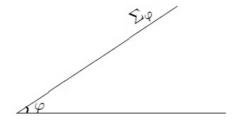


Fig. 1. The angle Σ_{φ} of degree $\varphi \in (0,\pi]$

The quadratic form in $L^2(\mathbb{R}^2)$

$$\mathfrak{a}_{\varphi}[f] := \|\nabla f\|_{L^{2}(\mathbb{R}^{2};\mathbb{C}^{2})}^{2} - \alpha \|f|_{\Sigma_{\varphi}}\|_{L^{2}(\Sigma_{\varphi})}^{2}, \qquad \operatorname{dom} \mathfrak{a}_{\varphi} := H^{1}(\mathbb{R}^{2}), \tag{1}$$

is closed, densely defined, symmetric and lower-semibounded, where $f|_{\Sigma_{\varphi}}$ is the trace of f on Σ_{φ} , and the constant $\alpha>0$ is called the strength of interaction. The corresponding self-adjoint operator in $L^2(\mathbb{R}^2)$ we denote by A_{φ} . Known spectral properties of this operator include explicit representation of the essential spectrum $\sigma_{\rm ess}(A_{\varphi})=[-\alpha^2/4,+\infty)$ and some

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information on the discrete spectrum: $\sharp \sigma_{\rm d}(A_\varphi) \geqslant 1$ if and only if $\varphi \neq \pi$. These two statements can be deduced from more general results by Exner and Ichinose [EI01]. They are complemented by Exner and Nemčová [EN03] with the limiting property $\sharp \sigma_{\rm d}(A_\varphi) \to +\infty$ as $\varphi \to 0+$.

In [LLP10] the author obtained jointly with Igor Lobanov and Igor Yu. Popov a general result, which implies the lower bound on the spectrum of A_{φ}

$$\inf \sigma(A_{\varphi}) \geqslant -\frac{\alpha^2}{4\sin^2(\varphi/2)}.$$
 (2)

This bound is close to optimal for φ close to π , whereas in the limit $\varphi \to 0+$ the bound tends to $-\infty$. In the present note we sharpen this bound. Namely, we obtain

$$\inf \sigma(A_{\varphi}) \geqslant -\frac{\alpha^2}{(1+\sin(\varphi/2))^2}.$$
 (3)

The new bound yields that the operators A_{φ} are uniformly lower-semibounded with respect to φ and

$$\inf \sigma(A_{\varphi}) \geqslant -\alpha^2$$

holds for all $\varphi \in (0, \pi]$. This observation agrees well with physical expectations. Note that separation of variables yields that $\inf \sigma(A_{\pi}) = -\alpha^2/4$ and in this case the lower bound in (3) coincides with the exact spectral bottom.

For sufficiently sharp angles upper bounds on $\inf \sigma(A_{\varphi})$ were obtained by Brown, Eastham and Wood [BEW08]. See also Open Problem 7.3 [E08] related to the discrete spectrum of A_{φ} for φ close to π .

δ-interactions on crossing straight lines. We also consider an analogous model with the δ-interaction supported on the set $\Gamma_{\varphi} = \Gamma_1 \cup \Gamma_2$, where Γ_1 and Γ_2 are two straight lines, which cross at the angle $\varphi \in (0, \pi)$ as in Figure 2.

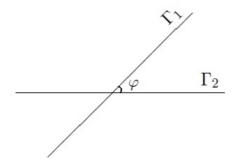


Fig. 2. The straight lines Γ_1 and Γ_2 crossing at the angle of degree $\varphi \in (0,\pi)$

The corresponding self-adjoint operator B_{φ} in $L^2(\mathbb{R}^2)$ can be defined via the closed, densely defined, symmetric and lower-semibounded quadratic form in $L^2(\mathbb{R}^2)$

$$\mathfrak{b}_{\varphi}[f] := \|\nabla f\|_{L^{2}(\mathbb{R}^{2};\mathbb{C}^{2})}^{2} - \alpha \|f|_{\Gamma_{\varphi}}\|_{L^{2}(\Gamma_{\varphi})}^{2}, \qquad \operatorname{dom} \mathfrak{b}_{\varphi} := H^{1}(\mathbb{R}^{2}), \tag{4}$$

where $\alpha>0$ is the strength of interaction. According to [EN03] it is known that $\sigma_{\rm ess}(B_\varphi)=[-\alpha^2/4,+\infty)$ and that $\sharp\sigma_{\rm d}(B_\varphi)\geqslant 1$.

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In this note we obtain the lower bound

$$\inf \sigma(B_{\varphi}) \geqslant -\frac{\alpha^2}{1 + \sin \varphi},\tag{5}$$

using the same method as for the operator A_{φ} . Separation of variables yields $\inf \sigma(B_{\pi/2}) = -\alpha^2/2$, and in this case the lower bound in the estimate (5) coincides with the exact spectral bottom.

Upper bounds on $\inf \sigma(B_{\varphi})$ were obtained in [BEW08, BEW09]. The operators of the type B_{φ} arise in the one-dimensional quantum three-body problem after excluding the center of mass, see Cornean, Duclos and Ricaud [CDR06, CDR08] and the references therein.

We want to stress that our proofs are of elementary nature and we do not use any reduction to integral operators acting on interaction supports Σ_{φ} and Γ_{φ} .

2. Sobolev spaces on wedges

In this section $\Omega \subset \mathbb{R}^2$ is a wedge with an angle of $\varphi \in (0, 2\pi)$. The Sobolev space $H^1(\Omega)$ is defined as usual, see [McL, Chapter 3]. For any $f \in H^1(\Omega)$ the trace $f|_{\partial\Omega} \in L^2(\partial\Omega)$ is well-defined as in [McL, Chapter 3] and [M87].

Proposition 2.1. [LP08, Lemma 2.6] Let Ω be a wedge with angle of degree $\varphi \in (0, \pi]$. Then for any $f \in H^1(\Omega)$ the estimate

$$\|\nabla f\|_{L^{2}(\Omega;\mathbb{C}^{2})}^{2} - \gamma \|f|_{\partial\Omega}\|_{L^{2}(\partial\Omega)}^{2} \geqslant -\frac{\gamma^{2}}{\sin^{2}(\varphi/2)} \|f\|_{L^{2}(\Omega)}^{2}$$

holds for all $\gamma > 0$.

Proposition 2.2. [LP08, Lemma 2.8] Let Ω be a wedge with angle of degree $\varphi \in (\pi, 2\pi)$. Then for any $f \in H^1(\Omega)$ the estimate

$$\|\nabla f\|_{L^{2}(\Omega;\mathbb{C}^{2})}^{2} - \gamma \|f|_{\partial\Omega}\|_{L^{2}(\partial\Omega)}^{2} \geqslant -\gamma^{2} \|f\|_{L^{2}(\Omega)}^{2}$$

holds for all $\gamma > 0$.

Propositions 2.1 and 2.2 are variational equivalents of spectral results from [LP08].

3. A lower bound on the spectrum of A_{φ}

In the next theorem we sharpen the bound (2) using only properties of the Sobolev space H^1 on wedges and some optimization.

Theorem 3.1. Let the self-adjoint operator A_{φ} be associated with the quadratic form given in (1). Then the estimate

$$\inf \sigma(A_{\varphi}) \geqslant -\frac{\alpha^2}{\left(1 + \sin(\varphi/2)\right)^2}$$

holds.

Proof. The angle Σ_{φ} separates the Euclidean space \mathbb{R}^2 into two wedges Ω_1 and Ω_2 with angles of degrees φ and $2\pi - \varphi$, see Figure 3.

The underlying Hilbert space can be decomposed as

$$L^2(\mathbb{R}^2) = L^2(\Omega_1) \oplus L^2(\Omega_2).$$

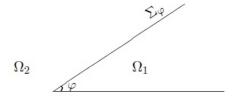


Fig. 3. The angle Σ_φ separates the Euclidean space \mathbb{R}^2 into two wedges Ω_1 and Ω_2

Any $f \in \text{dom } \mathfrak{a}_{\varphi}$ can be written as the orthogonal sum $f_1 \oplus f_2$ with respect to that decomposition of $L^2(\mathbb{R}^2)$. Note that $f_1 \in H^1(\Omega_1)$ and that $f_2 \in H^1(\Omega_2)$. Clearly,

$$||f||_{L^{2}(\mathbb{R}^{2})}^{2} = ||f_{1}||_{L^{2}(\Omega_{1})}^{2} + ||f_{2}||_{L^{2}(\Omega_{2})}^{2}, ||\nabla f||_{L^{2}(\mathbb{R}^{2};\mathbb{C}^{2})}^{2} = ||\nabla f_{1}||_{L^{2}(\Omega_{1};\mathbb{C}^{2})}^{2} + ||\nabla f_{2}||_{L^{2}(\Omega_{2};\mathbb{C}^{2})}^{2}.$$

$$(6)$$

The coupling constant can be decomposed as $\alpha = \beta + (\alpha - \beta)$ with some optimization parameter $\beta \in [0, \alpha]$ and the relation

$$\alpha \|f|_{\Sigma_{\varphi}}\|_{L^{2}(\Sigma_{\varphi})}^{2} = \beta \|f_{1}|_{\partial\Omega_{1}}\|_{L^{2}(\partial\Omega_{1})}^{2} + (\alpha - \beta) \|f_{2}|_{\partial\Omega_{2}}\|_{L^{2}(\partial\Omega_{2})}^{2}.$$
 (7)

holds. According to Proposition 2.1

$$\|\nabla f_1\|_{L^2(\Omega_1;\mathbb{C}^2)}^2 - \beta \|f_1|_{\partial\Omega_1}\|_{L^2(\partial\Omega_1)}^2 \geqslant -\frac{\beta^2}{\sin^2(\varphi/2)} \|f_1\|_{L^2(\Omega_1)}^2, \tag{8}$$

and according to Proposition 2.2

$$\|\nabla f_2\|_{L^2(\Omega_2;\mathbb{C}^2)}^2 - (\alpha - \beta)\|f_2|_{\partial\Omega_2}\|_{L^2(\partial\Omega_2)}^2 \geqslant -(\alpha - \beta)^2\|f_2\|_{L^2(\Omega_2)}^2. \tag{9}$$

The observations (6), (7) and the estimates (8), (9) imply

$$\mathfrak{a}_{\varphi}[f] \geqslant -\max\left\{\frac{\beta^2}{\sin^2(\varphi/2)}, (\alpha-\beta)^2\right\} \|f\|_{L^2(\mathbb{R}^2)}^2.$$

Making optimization with respect to β , we observe that the maximum between the two values in the estimate above is minimal, when these two values coincide. That is

$$\frac{\beta^2}{\sin^2(\varphi/2)} = (\alpha - \beta)^2,$$

which is equivalent to

$$\beta = \frac{\alpha \sin(\varphi/2)}{(1 + \sin(\varphi/2))},\tag{10}$$

resulting in the final estimate

$$\mathfrak{a}_{\varphi}[f] \geqslant -\frac{\alpha^2}{(1+\sin(\varphi/2))^2} ||f||_{L^2(\mathbb{R}^2)}^2.$$

This final estimate implies the desired spectral bound.

Remark 3.2. Note that the previously known lower bound (2) was derived from the proof of the last theorem if we choose $\beta = \alpha/2$, which is the optimal choice in our proof only for $\varphi = \pi$ as we see from (10).

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4. A lower bound on the spectrum of \mathbf{B}_{φ}

In the next theorem we obtain a lower bound on the spectrum of the self-adjoint operator B_{φ} using the same idea as in Theorem 3.1.

Theorem 4.1. Let the self-adjoint operator B_{φ} be associated with the quadratic form given in (4). Then the estimate

$$\inf \sigma(B_{\varphi}) \geqslant -\frac{\alpha^2}{1 + \sin \varphi}$$

holds.

Proof. The crossing straight lines Γ_1 and Γ_2 separate the Euclidean space \mathbb{R}^2 into four wedges $\{\Omega_k\}_{k=1}^4$. Namely, the wedges Ω_1 and Ω_2 with angles of degree φ and the wedges Ω_3 and Ω_4 with angles of degree $\pi - \varphi$, see Figure 4.

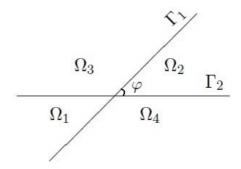


Fig. 4. The crossing straight lines Γ_1 and Γ_2 separate the Euclidean space \mathbb{R}^2 into four wedges $\{\Omega_k\}_{k=1}^4$

The underlying Hilbert space can be decomposed as

$$L^2(\mathbb{R}^2) = \bigoplus_{k=1}^4 L^2(\Omega_k).$$

Any $f \in \text{dom } \mathfrak{b}_{\varphi}$ can be written as the orthogonal sum $\bigoplus_{k=1}^{4} f_{k}$ with respect to that decomposition of $L^{2}(\mathbb{R}^{2})$. Note that $f_{k} \in H^{1}(\Omega_{k})$ for k = 1, 2, 3, 4. Clearly,

$$||f||_{L^{2}(\mathbb{R}^{2})}^{2} = \sum_{k=1}^{4} ||f_{k}||_{L^{2}(\Omega_{k})}^{2}, \quad ||\nabla f||_{L^{2}(\mathbb{R}^{2};\mathbb{C}^{2})}^{2} = \sum_{k=1}^{4} ||\nabla f_{k}||_{L^{2}(\Omega_{k};\mathbb{C}^{2})}^{2}.$$
(11)

The coupling constant can be decomposed as $\alpha = \beta + (\alpha - \beta)$ with some optimization parameter $\beta \in [0, \alpha]$ and the relation

$$\alpha \|f|_{\Gamma_{\varphi}}\|_{L^{2}(\Gamma_{\varphi})}^{2} = \beta \|f_{1}|_{\partial\Omega_{1}}\|_{L^{2}(\partial\Omega_{1})}^{2} + \beta \|f_{2}|_{\partial\Omega_{2}}\|_{L^{2}(\partial\Omega_{2})}^{2} + (\alpha - \beta) \|f_{3}|_{\partial\Omega_{3}}\|_{L^{2}(\partial\Omega_{3})}^{2} + (\alpha - \beta) \|f_{4}|_{\partial\Omega_{4}}\|_{L^{2}(\partial\Omega_{4})}^{2}$$
(12)

holds. According to Proposition 2.1

$$\|\nabla f_1\|_{L^2(\Omega_1;\mathbb{C}^2)}^2 - \beta \|f_1|_{\partial\Omega_1}\|_{L^2(\partial\Omega_1)}^2 \geqslant -\frac{\beta^2}{\sin^2(\varphi/2)} \|f_1\|_{L^2(\Omega_1)}^2,$$

$$\|\nabla f_2\|_{L^2(\Omega_2;\mathbb{C}^2)}^2 - \beta \|f_2|_{\partial\Omega_2}\|_{L^2(\partial\Omega_2)}^2 \geqslant -\frac{\beta^2}{\sin^2(\varphi/2)} \|f_2\|_{L^2(\Omega_2)}^2.$$
(13)

Also according to Proposition 2.1

$$\|\nabla f_3\|_{L^2(\Omega_3;\mathbb{C}^2)}^2 - (\alpha - \beta)\|f_3|_{\partial\Omega_3}\|_{L^2(\partial\Omega_3)}^2 \geqslant -\frac{(\alpha - \beta)^2}{\cos^2(\varphi/2)}\|f_3\|_{L^2(\Omega_3)}^2,$$

$$\|\nabla f_4\|_{L^2(\Omega_4;\mathbb{C}^2)}^2 - (\alpha - \beta)\|f_4|_{\partial\Omega_4}\|_{L^2(\partial\Omega_4)}^2 \geqslant -\frac{(\alpha - \beta)^2}{\cos^2(\varphi/2)}\|f_4\|_{L^2(\Omega_4)}^2.$$
 (14)

The observations (11), (12) and the estimates (13), (14) imply

$$\mathfrak{b}_{\varphi}[f] \geqslant -\max\left\{\frac{\beta^2}{\sin^2(\varphi/2)}, \frac{(\alpha-\beta)^2}{\cos^2(\varphi/2)}\right\} \|f\|_{L^2(\mathbb{R}^2)}^2.$$

Making optimization with respect to β , we observe that the maximum between the two values in the estimate above is minimal, when these two values coincide. That is

$$\frac{\beta^2}{\sin^2(\varphi/2)} = \frac{(\alpha - \beta)^2}{\cos^2(\varphi/2)},$$

which is equivalent to

$$\beta = \frac{\alpha \tan(\varphi/2)}{(1+\tan(\varphi/2))},\tag{15}$$

resulting in the final estimate

$$\mathfrak{b}_{\varphi}[f] \geqslant -\frac{\alpha^2}{1+\sin(\varphi)} ||f||_{L^2(\mathbb{R}^2)}^2.$$

This final estimate implies the desired spectral bound.

Remark 4.2. The result of Theorem 4.1 complements [CDR08, Theorem 4.6 (iv)], where the bound

$$\inf \sigma(B_{\varphi}) \geqslant -\alpha^2$$
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for all $\varphi \in (0, \pi)$ was obtained.

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