Self-similar solutions of a cross-diffusion parabolic system with variable density: explicit estimates and asymptotic behaviour

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PACS 02.30.Jr, 02.30.Mv, 11.10.Jj, 11.10.Lm

DOI 10.17586/2220-8054-2017-8-1-5-12

In this paper, we study the properties of self-similar solutions of a cross-diffusion parabolic system. In particular, we find the Zeldovich– Barenblatt type solution to the cross diffusive system. The asymptotic behavior of self-similar solutions are analyzed for both the slow and fast diffusive regimes. It is shown that coefficients of the main term of the asymptotic of solution satisfy some system of nonlinear algebraic equations.

Keywords: cross-diffusive system, non-divergence form, finite speed, perturbation, global solutions, asymptotic behavior, numerical analysis.

Received: 25 July 2016

Revised: 28 August 2016

1. Introduction

Consider in $Q = \{(t, x) : t > 0, x \in \mathbb{R}^N\}$ the cross-diffusive system of equations in non-divergence form with Cauchy conditions:

$$|x|^{n} \frac{\partial u}{\partial t} = v^{\alpha_{1}} \nabla \left(|x|^{k} u^{m_{1}-1} \nabla u \right) + |x|^{n} u^{\beta_{1}},$$

$$|x|^{n} \frac{\partial v}{\partial t} = u^{\alpha_{2}} \nabla \left(|x|^{k} v^{m_{1}-1} \nabla v \right) + |x|^{n} v^{\beta_{2}},$$
(1)

$$u(0,x) = u_0(x), \quad v(0,x) = v_0(x), x \in \mathbb{R}^N.$$
 (2)

where $n, k, m_1, m_2, \alpha_1, \alpha_2, \beta_1, \beta_2$ are the numerical parameters, $\nabla(\cdot) = \operatorname{grad}_x(\cdot)$, are the functions $u = u(t, x) \ge 0$, $v = v(t, x) \ge 0$ are the solutions.

It is clear that the system (1) is degenerate. Therefore, it does not have classical solutions on the domain defined by equations u(t,x) = 0, v(t,x) = 0, $\nabla u = 0$, $\nabla v = 0$, meaning system (1) may not have a classical solution. Therefore, in this case we consider a weak solution having the property $u(t,x) \ge 0$, $v(t,x) \ge 0$; $v^{\alpha_1} \nabla \left(|x|^k u^{m_1-1} \nabla u \right)$, $u^{\alpha_2} \nabla \left(|x|^k v^{m_1-1} \nabla v \right) \in C(Q)$ and obeying to the system (1) in sense of a distribution [1].

Non-divergent form equations and system of equations (1) are often used to describe various physical phenomena, such as the diffusive process for biological species, the resistive diffusion phenomena in force-free magnetic fields, curve shortening flow, spreading of infectious disease and so on, see for [2–8].

In [2] for the system (1) (for n = 0, k = 0, and $m_1 = 1$, $m_2 = 1$) the authors proved that: (i) when $\min \{a, b\} \le \lambda_1$, then there exists global positive classical solution, and all positive classical solutions cannot blow up in finite time in the meaning of maximum norm; (ii) when $\min \{a, b\} > \lambda_1$, there is no global positive classical solution.

Previous studies [3] with positive solutions of some degenerate and quasilinear parabolic systems are given by:

$$u_{it} = f_i (u_{i+1}) (\Delta u_i + a_i u_i)$$

$$u_{nt} = f_n (u_1) (\Delta u_n + a_n u_n), \quad x \in \Omega, \ t > 0, \ i = 1, 2, ..., n - 1.$$

The local existence and uniqueness of classical solution also proven in [3]. It is also shown that: (i) when $\min \{a_1, ..., a_n\} \leq \lambda_1$ then there exists global positive classical solution, and all positive classical solutions cannot blow up in finite time in the meaning of maximum norm; (ii) when $\min \{a_1, ..., a_n\} > \lambda_1$, and the initial datum $(u_{10}, ..., u_{n0})$ satisfies some assumptions, then the positive classical solution is unique and blows up in finite time, where λ_1 is the first eigenvalue of Δ in Ω with homogeneous Dirichlet boundary conditions.

Some properties of solutions for the non-divergence form in single equation, such as the existence, nonuniqueness, the blow-up properties etc. have been discussed by many authors [4–6]. Finite speed properties of a perturbation of distribution (FSPD) and the asymptotic behavior of self-similar solutions for another systems are considered in [9–12].

In [13,14], the Cauchy problem for the following two equations with variable coefficients is studied:

$$\rho(x)\frac{\partial u}{\partial t} = \operatorname{div}\left(u^{m-1}|\nabla u|^{p-2}\nabla u\right) + \rho(x)u^{\beta},$$

and

$$\rho(x)\frac{\partial u}{\partial t} = \operatorname{div}\left(u^{m-1}|\nabla u|^{p-2}\nabla u\right) + u^{\beta}, \ x \in \mathbb{R}^{N}, \ t > 0,$$

where p > 1, m + p - 3 > 0, $\beta > m + p - 2$, $\rho(x) = |x|^{-n}$ or $\rho(x) = (1 + |x|)^{-n}$. The authors showed that under some constraints for the parameters and for the initial data, any nontrivial solution of the Cauchy problem blows up in finite time. Moreover, the authors established a sharp universal estimate of the solution near the blow-up point.

Ref. [7] considers the degenerate and quasilinear parabolic systems having non divergent form given by:

$$\begin{cases} u_{it} = u_i^{p_i} \left(\Delta u_i + a_i u_{i+1} \right), \ i = 1, 2, ..., m, \ u_{m+1} = u_1 \\ u_i \left(x, 0 \right) = u_{i0} \left(x \right), \ i = 1, 2, ..., m, \ x \in \Omega, \\ u_i \left(x, t \right) = 0, \ i = 1, 2, ..., m, \ x \in \partial\Omega, \ t > 0, \end{cases}$$

where a_i , p_i are the positive constants, $\Omega \in \mathbb{R}^N$ is a bounded domain with smooth boundary $\partial \Omega$. The local existence and uniqueness of classical solution are proved. Moreover, it proved that all solutions exist globally with homogeneous Dirichlet boundary conditions.

Ref. [8], investigates positive solutions of the degenerate parabolic equation not in divergence form: $u_t = u^p \Delta u + u^q |\nabla u|^2 - u^r$, subject to null Dirichlet boundary conditions. The existence of global solutions and the large time behavior for them is studied. The main effort is paid to obtain uniform asymptotic profiles for decay solutions, under various dominations of the nonlinear diffusion or absorption. It is shown that the large time property of the solution u behaves just like (1 + (r - 1)t)(-1/r - 1) if the decay is governed by the nonlinear absorption with 1 < r < p + 1.

In this paper we consider the problem of constructing Zeldovich–Barenblatt type solution for the cross system equation (1). Using comparison methods, the property of FSPD of the Cauchy problem for a cross-diffusion parabolic system not in divergence form is established. The asymptotic behavior is discussed for a solution of the cross-diffusion parabolic system equations in non-divergence form for slow and fast diffusion cases depending on the value of the numerical parameters. On the basis of the asymptotic of solutions, suitable initial approximations are offered for the iterative process in the cases of the slow and fast diffusions, depending on the numerical parameter values.

2. The self-similar system of equations

Here, we provide a method of nonlinear splitting [15] for construction of self-similar equation for the system given by Eqs. (1). We look for the solutions u(t, x), v(t, x) in the form:

$$u(t,x) = \bar{u}(t) w(\tau(t),r), v(t,x) = \bar{v}(t) \psi(\tau(t),r).$$
(3)

Then, we obtain $\bar{u}(t)$, $\bar{v}(t)$ as $\bar{u}(t) = (T+t)^{1/(1-\beta_1)}$, $\bar{v}(t) = (T+t)^{1/(1-\beta_2)}$. From Eqs. (3) and (1), we obtain the following system of equations:

$$r^{n}\frac{\partial w}{\partial \tau} = \psi^{\alpha_{1}}r^{1-N}\frac{\partial}{\partial r}\left(r^{N-1+k}w^{m_{1}-1}\frac{\partial w}{\partial r}\right) + b_{1}r^{n}\tau^{-1}\left(w^{\beta_{1}} - \frac{1}{1-\beta_{1}}w\right),$$

$$r^{n}\frac{\partial \psi}{\partial \tau} = w^{\alpha_{2}}r^{1-N}\frac{\partial}{\partial r}\left(r^{N-1+k}\psi^{m_{2}-1}\frac{\partial \psi}{\partial r}\right) + b_{2}r^{n}\tau^{-1}\left(\psi^{\beta_{2}} - \frac{1}{1-\beta_{2}}\psi\right),$$
(4)

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where:

$$\begin{split} \tau\left(t\right) &= \int \left(T+t\right)^{\frac{\alpha_1}{1-\beta_2} + \frac{m_1-1}{1-\beta_1}} dt = \int \left(T+t\right)^{\frac{\alpha_2}{1-\beta_1} + \frac{m_2-1}{1-\beta_2}} dt \quad \text{at} \quad \frac{\alpha_1}{1-\beta_2} + \frac{m_1-1}{1-\beta_1} + 1 \neq 0, \\ \tau\left(t\right) &= \ln\left(T+t\right) \quad \text{at} \quad \frac{\alpha_1}{1-\beta_2} + \frac{m_1-1}{1-\beta_1} + 1 = 0, \\ b_i &= \left(\frac{\alpha_i}{1-\beta_{3-i}} + \frac{m_i-1}{1-\beta_i} + 1\right)^{-1}, \ i = 1, 2, \ r = |x|, \\ \frac{\alpha_1}{1-\beta_2} + \frac{m_1-1}{1-\beta_1} = \frac{\alpha_2}{1-\beta_1} + \frac{m_2-1}{1-\beta_2}. \end{split}$$

It is easy to show that the system (4) has approximately self-similar solution given by:

$$w(\tau, r) = f(\xi), \quad \psi(\tau, r) = \phi(\xi), \quad \xi = r^{\frac{2-k}{2}} / \tau^{\frac{2-k}{4-2k+2n}}, \tag{5}$$

where ξ is self-similar variable and the functions $f(\xi)$, $\phi(\xi)$ fulfill the following approximately self-similar system of equations:

$$\phi^{\alpha_1}\xi^{1-\frac{2N}{2-k}}\frac{d}{d\xi}\left(\xi^{\frac{2N}{2-k}-1}f^{m_1-1}\frac{df}{d\xi}\right) + \frac{2}{(2-k)(n+2-k)}\xi^{\frac{2N}{2-k}+1}\frac{df}{d\xi} + \frac{4b_1}{(2-k)^2}\xi^{\frac{2N}{2-k}}\left(f^{\beta_1} - \frac{1}{1-\beta_1}f\right) = 0,$$

$$f^{\alpha_2}\xi^{1-\frac{2N}{2-k}}\frac{d}{d\xi}\left(\xi^{\frac{2N}{2-k}-1}\phi^{m_2-1}\frac{d\phi}{d\xi}\right) + \frac{2}{(2-k)(n+2-k)}\xi^{\frac{2N}{2-k}+1}\frac{d\phi}{d\xi} + \frac{4b_2}{(2-k)^2}\xi^{\frac{2N}{2-k}}\left(\phi^{\beta_2} - \frac{1}{1-\beta_2}\phi\right) = 0.$$
(6)

In the following, we will consider nontrivial, nonnegative solutions of the system (6) satisfying the following conditions:

$$f(0) = M_1, \ \phi(0) = M_2, \ M_1 \in \mathbb{R}^+, M_2 \in \mathbb{R}^+, f(d_1) = \phi(d_2) = 0, \ 0 < d_1 < \infty, \ 0 < d_2 < \infty.$$

$$(7)$$

3. The slow diffusion case $(m_{3-i} > 1 + \alpha_i, i = 1, 2)$. Explicit estimates of the solution of problem (1)-(2)

Using the solution comparison method of [1] and the standard equations method of [15] for solving the problem (6)–(7), we can obtain the estimates for the solution of the problem (1)–(2).

We note that the functions:

$$\overline{f}(\xi) = A \left(a - \xi^{\frac{2n}{2-k}+2} \right)_{+}^{p_1}, \quad \overline{\phi}(\xi) = B \left(a - \xi^{\frac{2n}{2-k}+2} \right)_{+}^{p_2}, \tag{8}$$

where $p_i = \frac{m_{3-i} - \alpha_i - 1}{(m_1 - 1)(m_2 - 1) - \alpha_1 \alpha_2}$, $i = 1, 2, (b)_+ = \max(0, b)$, $a^{p_1}A = M_1$, $a^{p_2}B = M_2$ at $\xi < a^{\left(\frac{2n}{2-k}+2\right)^{-1}}$ satisfy Eq. (7).

Theorem 1. Let $m_{3-i} > 1 + \alpha_i$, $p_i m_i - 1 > 0$, i = 1, 2,

$$\begin{split} M_{i}^{\beta_{i}-1} &\leq \frac{1}{1-\beta_{i}} + \frac{N+n}{n+2-k} \frac{p_{i}}{b_{i} \left(p_{i}m_{i}-1\right)}, \quad i = 1, 2, \\ M_{i}^{m_{i}-1}M_{3-i}^{\alpha_{i}} &= \frac{a}{\left(p_{i}m_{i}-1\right)\left(n+2-k\right)^{2}}, \quad i = 1, 2, \\ u_{+}\left(0, x\right) &\geq u_{0}\left(x\right), \quad v_{+}\left(0, x\right) \geq v_{0}\left(x\right), \quad x \in \mathbb{R}^{N}. \end{split}$$

Then, in Q the problem (1)–(2) has a global solution with the estimate:

$$u(t,x) \leq u_{+}(t,x) = \bar{u}(t) f(\xi),$$

$$v(t,x) \leq v_{+}(t,x) = \bar{v}(t) \overline{\phi}(\xi).$$
(9)

Proof. Theorem 1 is proved by the method of comparison of solutions. As comparable functions, we take the functions $u_{+}(t, x)$, $v_{+}(t, x)$, defined by Eq. (9). Then, according to (3), (5) and $M_{i}^{m_{i}-1}M_{3-i}^{\alpha_{i}} = \frac{a}{(p_{i}m_{i}-1)(n+2-k)^{2}}$, i = 1, 2, we get:

$$Lu_{+}(t,x) = \left(-\frac{4p_{1}(N+n)}{(2-k)^{2}(n+2-k)(p_{1}m_{1}-1)} - \frac{4b_{1}}{(1-\beta_{1})(2-k)^{2}}\right)\xi^{\frac{2n}{2-k}}\overline{f}(\xi) + \frac{4b_{1}}{(2-k)^{2}}\xi^{\frac{2n}{2-k}}\overline{f}^{\beta_{1}}(\xi),$$

$$Lv_{+}(t,x) = \left(-\frac{4p_{2}(N+n)}{(2-k)^{2}(n+2-k)(p_{2}m_{2}-1)} - \frac{4b_{2}}{(1-\beta_{2})(2-k)^{2}}\right)\xi^{\frac{2n}{2-k}}\overline{\phi}(\xi) + \frac{4b_{2}}{(2-k)^{2}}\xi^{\frac{2n}{2-k}}\overline{\phi}^{\beta_{2}}(\xi).$$

From these expressions and $\overline{f}(\xi) \leq Aa^{p_1} = M_1, \ \overline{\phi}(\xi) \leq Ba^{p_2} = M_2$, we find, that for execution, the inequalities $Lu_{+} \leq 0$, $Lv_{+} \leq 0$ are enough to fulfill the conditions

$$M_i^{\beta_i - 1} \le \frac{1}{1 - \beta_i} + \frac{N + n}{n + 2 - k} \frac{p_i}{b_i (p_i m_i - 1)}, \quad i = 1, 2.$$

By the conditions of the theorem, the last inequalities are satisfied. Then, according to the theorem, for comparison of problem (1)-(2) solutions, there is a global solution in Q and the following estimation $u_+(t,x) \ge 0$ $u(t,x), v_+(t,x) \ge v(t,x), x \in \mathbb{R}^N$ holds.

We note for the solutions of Cauchy problem (1)-(2), we have FSPD properties.

The asymptotic behavior of self-similar solutions of the problem (6)-(7) 4.

Let us introduce the following notations:

$$b_{i1} = p_i (p_i m_i - 1), \ b_{i2} = -\frac{p_i}{(n+2-k)^2}, \ b_{i3} = \frac{b_i}{(n+2-k)^2 a}, \ i = 1, 2.$$

Assume $\frac{\alpha_1}{1-\beta_2} + \frac{m_1-1}{1-\beta_1} = \frac{\alpha_2}{1-\beta_1} + \frac{m_2-1}{1-\beta_2}$. Then the following theorem is valid: **Theorem 2.** Let $\frac{\alpha_1}{1-\beta_2} + \frac{m_1-1}{1-\beta_1} + 1 > 0$. Then compactly supported solution of the problem (6), (7) as $\xi \to a^{\left(\frac{2n}{2-k}+2\right)^{-1}}$ has the following asymptotic behavior:

$$f(\xi) = c_1 \left(a - \xi^{\frac{2n}{2-k}+2} \right)^{p_1} (1+o(1)),$$

$$\phi(\xi) = c_2 \left(a - \xi^{\frac{2n}{2-k}+2} \right)^{p_2} (1+o(1)),$$
(10)

if one of the the following conditions are fulfilled: 1 1

(1)
$$p_1 = \frac{1}{1 - \beta_1}$$
, $p_2 = \frac{1}{1 - \beta_2}$ and the coefficients c_i $(i = 1, 2)$ are the roots of the systems of nonlinear algebraic equations:
 $b_{i1}c_{i1}^{m_i-1}c_{i1}^{\alpha_i} + b_{i2} + b_{i2}c_{i1}^{\beta_i-1} = 0$, $i = 1, 2$

 $b_{i1}c_i^{m_i-1}c_{3-i}^{\alpha_i} + b_{i2} + b_{i3}c_i^{\beta_i-1} = 0, \quad i = 1, 2.$ (2) $p_i < \frac{1}{1-\beta_i}, \quad p_i m_i > 1, \quad i = 1, 2 \text{ and the coefficients } c_i \quad (i = 1, 2) \text{ are the roots of the systems of nonlinear between the s$ algebraic equations:

$$b_{i1}c_i^{m_i-1}c_{3-i}^{\alpha_i}+b_{i2}=0, \ i=1,2.$$

(3) $p_1 = \frac{1}{1-\beta_1}$, $p_2 < \frac{1}{1-\beta_2}$, $p_2m_2 > 1$ and the coefficients c_i (i = 1, 2) are the roots of the systems of nonlinear algebraic equations:

$$\begin{cases} b_{11}c_1^{m_1-1}c_2^{\alpha_1} + b_{12} + b_{13}c_1^{\beta_1-1} = 0, \\ b_{21}c_2^{m_2-1}c_1^{\alpha_2} + b_{22} = 0. \end{cases}$$

(4) $p_1 < \frac{1}{1-\beta_1}$, $p_2 = \frac{1}{1-\beta_2}$, $p_1m_1 > 1$ and the coefficients c_i (i = 1, 2) are the roots of the systems of nonlinear algebraic equations:

$$\begin{cases} b_{11}c_1^{m_1-1}c_2^{\alpha_1} + b_{12} = 0, \\ b_{21}c_2^{m_2-1}c_1^{\alpha_2} + b_{22} + b_{23}c_2^{\beta_2-1} = 0. \end{cases}$$

Proof. To simplify such auxiliary systems of equations, one can use the following transformations:

$$f(\xi) = \overline{f}(\xi)y_1(\eta), \ \phi(\xi) = \overline{\phi}(\xi)y_2(\eta), \ \eta = -\ln\left(a - \xi^{\frac{2n}{2-k}+2}\right),$$
(11)

where $\bar{f}(\xi) = \left(a - \xi^{\frac{2n}{2-k}+2}\right)^{p_1}, \ \bar{\phi}(\xi) = \left(a - \xi^{\frac{2n}{2-k}+2}\right)^{p_2}$ and $y_1(\eta), \ y_2(\eta)$ - the new functions.

Now, let us study the asymptotic behavior of solutions of Eqs. (6) and (7) at $\xi \to a^{\left(\frac{2n}{2-k}+2\right)^{-1}}$. After transformation (11), Eqs. (6) and (7) will take the following form:

$$y_{3-i}^{\alpha_{i}}\frac{d}{d\eta}\left(Ly_{i}\right)+a_{i1}\left(\eta\right)y_{3-i}^{\alpha_{i}}Ly_{i}+a_{i2}\left(\eta\right)\left(\frac{dy_{i}}{d\eta}+a_{i0}\left(\eta\right)y_{i}\right)+a_{i3}\left(\eta\right)y_{i}+a_{i4}\left(\eta\right)y_{i}^{\beta_{i}}=0,\ (i=1,2).$$
 (12)

Here.

$$Ly_{i} = y_{i}^{m_{i}-1} \left(\frac{dy_{i}}{d\eta} + a_{i0}(\eta) y_{i} \right), \quad a_{i0}(\eta) = -p_{i}, \quad a_{i1}(\eta) = \frac{n+N}{n+2-k} \frac{e^{-\eta}}{a-e^{-\eta}} - p_{i}m_{i} + 1,$$

$$a_{i2}(\eta) = \frac{1}{(n+2-k)^{2}}, \quad a_{i3}(\eta) = -\frac{b_{i}}{(n+2-k)^{2}(1-\beta_{i})} \frac{e^{-\eta}}{a-e^{-\eta}}, \quad a_{i4}(\eta) = \frac{b_{i}}{(n+2-k)^{2}} \frac{e^{-(1+p_{i}(\beta_{i}-1))\eta}}{a-e^{-\eta}},$$

$$i = 1, 2.$$

We assumed that $\xi \in [\xi_0, \xi_1), 0 < \xi_0 < \xi_1, \xi_1 = a^{\left(\frac{2n}{2-k}+2\right)^{-1}}$. Therefore, the function $\eta(\xi)$ has the following properties:

$$\eta'(\xi) > 0 \text{ at } \xi \in [\xi_0, \xi_1), \quad \eta_0 = \eta(\xi_0) > 0, \quad \lim_{\xi \to \xi_1} \eta(\xi) = +\infty.$$

The auxiliary system of equations (12) studied under the following limitations:

$$\lim_{\eta \to +\infty} a_{ij}(\eta) = a_{ij}^0 \quad (i = 1, 2; \ j = 0, 1, 2, 3, 4),$$

exist and are finite and nonzero, that is $0 < |a_{ij}^0| < +\infty$.

The solutions of the system (12), in a certain vicinity of $+\infty$ satisfy the inequalities:

$$y_i(\eta) > 0, \quad y'_i + a_{i0}(\eta)y_i \neq 0 \quad (i = 1, 2).$$

Assuming in Eqs. (12):

$$v_{i}(\eta) = y_{i}^{m_{i}-1} \left(\frac{dy_{i}}{d\eta} + a_{i0}(\eta) y_{i} \right), \quad i = 1, 2,$$
(13)

we obtain the identity:

$$v_{i}^{'}(\eta) \equiv -a_{i1}(\eta)v_{i}(\eta) - a_{i2}(\eta)v_{i}(\eta)y_{i}^{1-m_{i}}y_{3-i}^{-\alpha_{i}} - a_{i3}(\eta)y_{i}y_{3-i}^{-\alpha_{i}} - a_{i4}(\eta)y_{i}^{\beta_{i}}y_{3-i}^{-\alpha_{i}}, \ (i = 1, 2).$$
(14)

Furthermore, consider the functions:

$$g_i(\lambda_i,\eta) \equiv -a_{i1}(\eta)\lambda_i - a_{i2}(\eta)\lambda_i y_i^{1-m_i} y_{3-i}^{-\alpha_i} - a_{i3}(\eta)y_i y_{3-i}^{-\alpha_i} - a_{i4}(\eta)y_i^{\beta_i} y_{3-i}^{-\alpha_i},$$
(15)

where $\lambda_i \in \mathbb{R}$, (i = 1, 2). Let us suppose $\frac{\alpha_1}{1 - \beta_2} + \frac{m_1 - 1}{1 - \beta_1} + 1 > 0$. Then,

$$\lim_{\eta \to +\infty} a_{i1}(\eta) = -p_i m_i + 1, \quad \lim_{\eta \to +\infty} a_{i3}(\eta) = 0, \quad \lim_{\eta \to +\infty} a_{i2}(\eta) = \frac{1}{(n+2-k)^2}$$
$$\lim_{\eta \to +\infty} a_{i4}(\eta) = \begin{cases} \frac{b_i}{a(n+2-k)^2}, & \text{at } 1 + p_i \left(\beta_i - 1\right) = 0, \\ 0, & \text{at } 1 + p_i \left(\beta_i - 1\right) > 0, \quad i = 1, 2. \end{cases}$$

and the functions $g_i(\lambda_i, \eta)$ (i = 1, 2) preserve sign on some interval $[\eta_1, +\infty) \subset [\eta_0, +\infty)$ for every fixed value λ_i (i = 1, 2).

Therefore, the functions $g_i(\lambda_i, \eta)$ (i = 1, 2) for all $\eta \in [\eta_1, +\infty)$ satisfies one of the inequalities:

$$g_i(\lambda_i, \eta) > 0 \quad \text{or} \quad g_i(\lambda_i, \eta) < 0 \quad (i = 1, 2).$$
 (16)

Let us assume that for the functions $v_i(\eta)$ (i = 1, 2) the limit at $\eta \to +\infty$ does not exist. Consider the case, where one of the inequalities (16) are fulfilled. By force of the function's $v_i(\eta)$ (i = 1, 2) oscillation, the straight line $\bar{v}_i = \lambda_i$ (i = 1, 2) and its graph intersects an infinite number of times in the interval $[\eta_i, +\infty)$ (i = 1, 2). However, this is impossible, since the interval $[\eta_i, +\infty)$ (i = 1, 2) is justly one of the inequalities (16) and therefore, it follows from Eq. (15), that graph of the functions $v_i(\eta)$ (i = 1, 2) intersects the straight line $\bar{v}_i = \lambda_i$ (i = 1, 2)only once in the interval $[\eta_i, +\infty)$ (i = 1, 2). Therefore, for the functions $v_i(\eta)$ (i = 1, 2) the limit at $\eta \to +\infty$ exists.

Assuming the functions $v_i(\eta)$ (i = 1, 2) are defined in accordance with Eq. (13) and have a limit at $\eta \to +\infty$, one can show that $y'_i(\eta)$ (i = 1, 2) has a limit at $\eta \to +\infty$, which is equal to zero. Then,

$$v_{i}(\eta) = y_{i}^{m_{i}-1} \left(\frac{dy_{i}}{d\eta} + a_{i0}(\eta) y_{i} \right) = a_{i0}^{0} \left(y_{i}^{0} \right)^{m_{i}} + o(1), \ i = 1, 2 \text{ at } \eta \to +\infty$$

and by (14) derivative of functions $v_i(\eta)$ (i = 1, 2) has a limit at $\eta \to +\infty$, which obviously equals zero. Consequently,

$$\lim_{\eta \to +\infty} \left(a_{i1}(\eta) v_i(\eta) + a_{i2}(\eta) v_i(\eta) y_i^{1-m_i} y_{3-i}^{-\alpha_i} + a_{i3}(\eta) y_i y_{3-i}^{-\alpha_i} + a_{i4}(\eta) y_i^{\beta_i} y_{3-i}^{-\alpha_i} \right) = 0 \quad (i = 1, 2).$$

It is easy to see that the system (12) has a solution $(y_1(\eta), y_2(\eta))$ with a finite non-zero limit at $\eta \to +\infty$ necessary to comply with the conditions of Theorem 2. Then the compactly supported solution of the problem (6)–(7) has an asymptotic of the form (10) as $\xi \to a^{\left(\frac{2n}{2-k}+2\right)^{-1}}$.

Thus, the theorem is proved.

Corollary 1. A weak solution of the problem (1)–(2) has the following asymptotic form:

$$u_A(t,x) \approx c_1(T+t)^{\frac{1}{1-\beta_1}} \left(a - \left(\frac{|x|^{\frac{2-k}{2}}}{\tau^{\left(\frac{2-k}{2-k}+2\right)^{-1}}}\right)^{\frac{2n}{2-k}+2} \right)^{p_1},$$
$$v_A(t,x) \approx c_2(T+t)^{\frac{1}{1-\beta_2}} \left(a - \left(\frac{|x|^{\frac{2-k}{2}}}{\tau^{\left(\frac{2-k}{2-k}+2\right)^{-1}}}\right)^{\frac{2n}{2-k}+2} \right)^{p_2},$$

at $|x| \to a^{\frac{1}{2-k+n}} \tau^{\frac{1}{n+2-k}}$, where constants c_1, c_2, p_1, p_2 are defined above.

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5. Fast diffusion case $(p_i < 0, i = 1, 2)$. Explicit estimates of the solution of problem (1)-(2)

Suppose that for system of equations (6) the following conditions are fulfilled:

$$f'(0) = 0, \ \phi'(0) = 0, \ f(\infty) = 0, \ \phi(\infty) = 0.$$
 (17)

Then, the functions

$$\overline{f}(\xi) = A_1 \left(a + \xi^{\frac{2n}{2-k}+2} \right)^{p_1}, \quad \overline{\phi}(\xi) = A_2 \left(a + \xi^{\frac{2n}{2-k}+2} \right)^{p_2}, \quad p_i = \frac{m_{3-i} - 1 - \alpha_i}{(m_1 - 1)(m_2 - 1) - \alpha_1 \alpha_2}, \quad i = 1, 2,$$

Theorem 3. Let $p_i < 0, \ \beta_i < 1, \ \frac{p_i}{1 - p_i m_i} \frac{n + N}{n + 2 - k} - \frac{b_i}{1 - \beta_i} + b_i a^{p_i(\beta_i - 1)} \ge 0,$ $A_i^{m_i - 1} A_{3-i}^{\alpha_i} = \frac{1}{(1 - p_i m_i)(n + 2 - k)^2}, \ i = 1, 2, \ u_0(x) \ge u_-(0, x), \ v_0(x) \ge v_-(0, x), \ x \in \mathbb{R}^N.$ Then, for problem (1) (2), d

Then, for problem (1)–(2), there is a global solution in Q and for it, the following estimation is correct: $u(t,x) \ge u_{-}(t,x), v(t,x) \ge v_{-}(t,x), x \in \mathbb{R}$, where

$$u_{-}(t,x) = \bar{u}(t)\overline{f}(\xi), \quad v_{-}(t,x) = \bar{v}(t)\overline{\phi}(\xi).$$
(18)

Theorem 3 is proved by the solution comparison method similar to the proof of Theorem 1. Furthermore, we introduce the following notations:

$$b_{i4} = p_i \left(\frac{n+N}{n+2-k} + p_i m_i - 1\right), \quad b_{i5} = \frac{p_i}{(n+2-k)^2} - \frac{b_i}{1 - \beta_i (n+2-k)^2}, \quad i = 1, 2$$

Also, we assume $\frac{\alpha_1}{1-\beta_2} + \frac{m_1-1}{1-\beta_1} = \frac{\alpha_2}{1-\beta_1} + \frac{m_2-1}{1-\beta_2}$. Then the following theorem is valid: **Theorem 4.** Let $\frac{\alpha_1}{1-\beta_2} + \frac{m_1-1}{1-\beta_1} + 1 > 0$, $\beta_i > 1$, $p_i < 0$, i = 1, 2. Then, vanishing at infinity, the solution of the number $\alpha_i = 1, 2$. the problem (6), (17) at $\xi \to \infty$ has an asymptotic:

$$f(\xi) = c_3 \left(a + \xi^{\frac{2n}{2-k}+2} \right)^{p_1}, \quad \phi(\xi) = c_4 \left(a + \xi^{\frac{2n}{2-k}+2} \right)^{p_2}$$

where the coefficients c_i (i = 3, 4) are the roots of systems of the nonlinear algebraic equations:

$$b_{i4}c_{i+2}^{m_i-1}c_{5-i}^{\alpha_i}+b_{i5}=0, \ i=1,2$$

The proof of the Theorem 4 is similar to that of the Theorem 2. **Corollary 2.** A weak solution of the problem (1)-(2) has the following asymptotic:

$$u_A(t,x) \approx c_3(T+t)^{\frac{1}{1-\beta_1}} \left(a + \left(\frac{|x|^{\frac{2-k}{2}}}{\tau^{\left(\frac{2-k}{2-k}+2\right)^{-1}}}\right)^{\frac{2n}{2-k}+2} \right)^{p_1},$$
$$v_A(t,x) \approx c_4(T+t)^{\frac{1}{1-\beta_2}} \left(a + \left(\frac{|x|^{\frac{2-k}{2}}}{\tau^{\left(\frac{2-k}{2-k}+2\right)^{-1}}}\right)^{\frac{2n}{2-k}+2} \right)^{p_2},$$

at $|x| \tau^{-\frac{1}{n+2-k}} \to +\infty$, where the constants c_3 , c_4 , p_1 , p_2 are defined above.

6. Numerical analysis of solutions

We note that due to the non-uniqueness of solutions, many different cases arise in the numerical study of problem (1)–(2). It is important to choose a proper initial approximation that preserves its nonlinearity properties. On the basis of the above qualitative studies, we produced numerical calculations. The numerical results show quick convergence of the iterative process to the solution of the Cauchy problem (1)–(2), due to the successful choice of the initial approximation. Below, some numerical experiment results for different numerical parameter values are presented.

Results of the numerical experiments and graphs show that the self-similar solutions are very appropriate approximations. Fig. 1 shows a compactly supported solution of the problem (1)–(2). Figs. 2 and 3 show the properties of the solution for problem (1)–(2), vanishing at infinity.



FIG. 1. Slow diffusion case $p_i > 0$, i = 1, 2: n < 0, k < 0 and n > 0, k < 0



FIG. 2. Fast diffusion case $p_i < 0$, i = 1, 2: n - k < 0 and n - k = 0

7. Conclusions

The Zeldovich–Barenblatt type solution for the cross diffusive system is obtained. The finite speed properties of the Cauchy problem for a cross-diffusion parabolic system with a source are studied using comparison method. The asymptotic behavior of self-similar solutions both for slow and fast diffusion cases are analyzed. It is shown that the coefficients in the main term of the asymptotic of the solution satisfy a system of nonlinear algebraic equations.

Results of the computational experiments show, that the self-similar solutions are very appropriate approximations and the iterative method based on the Picard method is effective for the numerical solution of nonlinear problems and leads to nonlinear effects if we use as initial approximation the solutions of self similar equations constructed by the nonlinear splitting and by the standard equation methods [9, 10, 15].



FIG. 3. Fast diffusion case $p_i < 0$, i = 1, 2: n - k = 1 and n - k > 1

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