The role of non-autonomous phases in the formation and transformation of solidphase oxide systems

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PACS 05.70.Np; 82.65.+r

ABSTRACT It is shown that the nature of processes in oxide solid-phase systems, primarily in nanosized ones, is determined by the behavior of the substance in a non-autonomous state. The composition of non-autonomous phases, the temperature of transition of non-autonomous phases to a liquid-like state, the equilibrium (locally equilibrium) and metastable thickness of non-autonomous phases, and the viscosity of the liquid-like nonautonomous phase are considered as the main parameters of non-autonomous phases.

KEYWORDS oxides, nanosystems, non-autonomous phase, melting point of the non-autonomous phase.

ACKNOWLEDGEMENTS The work was supported by state assignment No. 1023032900322-9-1.4.3

FOR CITATION Almjasheva O.V. The role of non-autonomous phases in the formation and transformation of solid-phase oxide systems. *Nanosystems: Phys. Chem. Math.*, 2024, **15** (6), 755–767.

1. Introduction

According to the definition given in the works of Defay and Prigogine [1, 2], surface phases, as well as phase and grain boundaries in polycrystalline systems are considered in a number of works as non-autonomous phases [4–46]. Such a definition has a number of advantages due to its more general nature. In this regard, the present paper employs this particular terminology.

A large number of both original works and reviews have been and are devoted to the formation and transformation of solid-phase oxides [46–66]. It should be noted that in the majority of the works on this topic, which have become classic, the main attention is paid to the influence, which the diffusion of components in bulk (autonomous) phases has on the processes occurring in the systems based on solid phases [61–66]. However, recent decades have shown an increased interest in studying the features of phase formation processes in nanopowders [47–61]. The volume fraction of nonautonomous phases in nanocrystalline systems, including those based on oxide nanocrystals, can be very significant [9]. This, apparently, should be taken into account in experimental studies, when determining the mechanisms and kinetics of the processes of phase formation and transformation, and when constructing physicochemical models of chemical and structural transformations in nanocrystalline systems. By now, there are no works that would analyze studies on the influence of the composition, structure, and properties of the non-autonomous phases on processes in solid-phase systems and, in the first place, in nanocrystalline systems. Of particular interest is the study of phase formation and phase transformation processes in materials based on oxide systems, since this is apparently the class of materials most widely represented in nature and in technology [65–73].

In relation to the listed reasons, this work is aimed at analyzing the results available in the literature, which are related to the determination of the properties and behavior of non-autonomous phases in polycrystalline systems and their influence on the course of phase formation processes and structural transformations, mainly in oxide phases.

2. Non-autonomous phases in polycrystals: composition, structure, properties

Conclusions that the composition, structure and surface properties of solid substances differ from their bulk properties were made on the basis of empirical data already in the works of 19th century scientists [74–76]. The works of Gibbs [77] theoretically consider some features of the behavior and properties of the surface of substances (interphase regions) from a thermodynamic standpoint. A detailed thermodynamic description of the regions between bulk (autonomous) phases as non-autonomous phases was given in [1–3]. These works considered the non-autonomous phases as regions between the coexisting (bordering on each other) autonomous phases, having a finite thickness, composition and structure, and changing from one autonomous phase to another.

It was shown in [4, 78–80] that such a description of non-autonomous phases is equivalent from the thermodynamic point of view to such a consideration of surface phases, in which all their 'excessive' properties are attributed to some

two-dimensional surface, as was proposed in [1–3]. This statement is absolutely true when the volume fraction of nonautonomous phases is negligibly small compared to the volume fraction of autonomous phases, i.e. for the case of macrocrystalline systems. When the particle (crystallite) sizes decrease to nanometer values, the volume fraction of non-autonomous phases increases sharply [11, 15]. In this case, it is no longer possible to ignore the thickness of nonautonomous phases. It is especially important to know the sizes and behavioral features of non-autonomous phases when considering the kinetics of processes in solid-phase systems [15].

The work [81] shows a certain similarity in the behavior of substance at the boundaries of different autonomous phases and at the boundaries of grains of one phase. That is why further on, as is shown in [81], non-autonomous phases can be understood as layers of finite thickness also in the autonomous phases contact area and at the boundaries of grains of the same phase, if these layers have not formed their own autonomous phase.

A detailed analysis of experimental data and the construction of calculation models, carried out in [81], showed that in the case of high-angle boundaries of grains of the same autonomous phase of variable composition, components redistribution between the autonomous and non-autonomous phases can be described in the first approximation by a dependency shown in Fig. 1. This result was obtained assuming that the non-autonomous phases formed at high-angle boundaries of the autonomous phase grains contain a thin layer of amorphous substance (Fig. 2). Naturally, a dependency similar to that shown in Fig. 1 cannot be expected in cases with low-angle and special boundaries of grains.

FIG. 1. Dependence of the impurity component distribution between the autonomous and nonautonomous phases [81]

It should be noted that the described effect of components redistribution between the autonomous and non-autonomous phases in particles based on crystalline autonomous phases can lead to the formation of structures of the core-shell type. In these structures, the shell appears due to the segregation of some components located on the particles surface in the non-autonomous phase. This effect is most noticeable when obtaining nanocrystalline particles. To date, a large number of examples of such structures formation are known [82–87]. A variant of the mechanism for the core-shell nanoparticles formation is described in [87]. However, apparently, this is only a special case of such nanostructures formation, and this area will be expanded subsequently.

When different solid autonomous phases are in contact, then the composition of the corresponding non-autonomous phase at temperatures not much different from the eutectic temperature will be close to the eutectic composition [81]. It should be noted that this result allows for describing the phenomenon of eutectic melting from positions different from the traditional ones [88].

Another case where it is convenient to use the term 'non-autonomous phase', or more precisely, the 'non-autonomous state of substance', is the state of substance under conditions of spatial limitations, when the amount of substance is insufficient for the formation of an autonomous phase with a composition and structure characteristic of the specified temperature and pressure. Such states may be exemplified by nanoparticles of certain sizes, thin surface layers, as well as the substance in micropores, in nanochannels, and in the matrices of composites with inclusions of nanoparticles at a certain placement density [89–106]. In this case, when the spatial restrictions are chemically inert, a substance, being in a non-autonomous state, will be unable to change the initially specified composition, but will have a structure different from that of the autonomous phase of the same composition at the same temperature and pressure.

FIG. 2. Schematic representation of a 2-D non-autonomous phase (regions II, II', III) and relaxation processes in a polycrystalline system [81]. Regions: ν_i – inside autonomous (I, I'), structured (II,II') and amorphized (III) zones of the non-autonomous phase; ν_i – region coinciding with the amorphized zone of the 2-D non-autonomous phase; v_k – two grains in contact with each other together with the intergranular formation. **Processes:** A – establishment of the locally equilibrium state in regions ν_i by introducing corresponding changes in the structure and composition; B – establishment of the locally equilibrium thickness of region III without mutual turning of grains; C – turning and shift of grains until the crystal sublattices completely coincide – for polycrystals based on one autonomous phase, or until the energy of translational mismatch is minimized – for systems consisting of different phases

The possibility of a liquid phase appearance at the boundary of solids at a temperature below their melting point was exemplified by ice already in [74, 75], but the quantitative relationship between the melting temperatures of the surface and bulk of solids was revealed only after an of experimental series by Tammann et al. [107]. It should be noted that despite the establishment of a quantitative relationship between the solid phase melting point and the socalled Tammann temperature (T_T) in these works (Table 1), the discovered phenomenon was interpreted incorrectly. For instance, the authors of [107] believed that volume diffusion processes are activated at the Tammann temperature. Only after a series of similar experiments, it was concluded in [108] that the T_T value in processes described in [107] determines the temperature, above which mass transfer is activated at the boundaries of solid phase particles.

Materials	Tammann temperature	References
Metals	~ 0.33 T _m	[107]
Oxide	$\sim 0.57~\text{T}_m$	
Metals	~ 0.3 T _m	[108]

TABLE 1. Tammann temperature for different materials

Based on the analysis of [109], the work [110–112] presents the generalized data on the temperature dependency of the diffusion in the crystalline (autonomous) phase, on the surface, and at the boundaries of grains, i.e. in nonautonomous phases. Taking into account the results of [107,108], as well as the dependency of the surface phase thickness on temperature [11], these results can be illustrated by Fig. 3. The data presented in Fig. 3 allow concluding that the substance in the non-autonomous phase is found in the liquid (liquid-like) state at temperatures above T_T .

FIG. 3. Temperature dependences of the surface phase thickness (a) and the diffusion coefficients of the substance (b). For case (b), the solid lines are from [81], the dashed line is an extrapolation of the dependence presented in [81]

By now, a significant number of experimental and theoretical works analyze the transition of non-autonomous phases to the liquid (liquid-like) state [113–131]. A detailed thermodynamic analysis of the relationship between the melting temperatures of autonomous and non-autonomous phases was made in [81]. It was shown in [11] that in cases when non-autonomous phases contain a thin amorphous layer of substance, the dependency between the melting temperatures of the autonomous and non-autonomous phases can be represented by the expression

$$
T_{m,n} \approx \frac{1 - \alpha_H}{1 - \alpha_S} T_m,\tag{1}
$$

where $\alpha_H = \Delta H_{am}/\Delta H_m$; $\alpha_S = \Delta S_{am}/\Delta S_m$; T_m – is the autonomous phase melting temperature, K; $T_{m,n}$ – is the melting temperature of the non-autonomous phase of the same composition, K; ΔH_{am} – is the amorphization enthalpy; ΔH_m – is the melting enthalpy; ΔS_{am} – is the amorphization entropy; ΔS_m – is the melting entropy.

Further on, expression (1) was analyzed in [18] for a number of oxide substances, and the limits of variation of the non-autonomous phase melting temperature were determined as a function of the autonomous phase melting temperature for this class of compounds:

$$
T_{m,n} = \gamma T_m,\tag{2}
$$

where γ – is a dimensionless parameter that takes values from 0.55 to 0.75, but most often 0.65.

A comparison of expression (2) with that for the Tammann temperature given in Table 1 shows their good agreement. This result once again confirms that in physical sense the empirical dependency of the Tammann temperature on the substances melting point reflects the dependency of the temperature of the non-autonomous phase transition into the liquid (liquid-like) state on the melting point of the autonomous phase of the corresponding composition. It should be noted that, as was analyzed in [81], the temperature of the non-autonomous phase transition into the liquid-like state depends on the phase structure. For example, in cases of the coherent grain intergrowth, this temperature approaches the autonomous phase melting point [81].

Another important parameter for the analysis of high-temperature processes in polycrystalline systems is the thickness of non-autonomous phases. From the mid-20th century to the present, a large number of experimental, theoretical papers and reviews have been devoted to the problem of determining the thickness and structure of non-autonomous phases. The works $[2,115,116,122,126-132]$ can be cited as an example. An important component of these works is the determination of the dependency of the non-autonomous phase thickness on temperature. The most well-known expressions for the nonautonomous phase thickness as a function of temperature are presented in Table 2. Among the expressions given in this Table, it is possible to single out the dependency described in [81], which differs from the others by two factors. First, this dependency changes at the non-autonomous phase melting point, which was not reflected in other formulas (Table 2). Second, when approaching the melting point of the solid phase, its thickness does not tend to infinity, as in other expressions, but remains a finite (albeit large) value. In this case, the crystal sizes decrease, apparently, to nanometer values. Physically, this means, as was indicated in [81], that the melting of solids can be considered as the

autonomization of the liquid non-autonomous phase with the transition of the crystalline (autonomous) phase to nonautonomous nanoclusters in the melt. To confirm the possibility of such a consideration of the melting process, the work [81] cited [135–144], in which various experimental data and theoretical calculations allowed making conclusions that solid (crystalline) clusters are observed in the melt near the melting point. The existence of nanosized solid clusters above the melting point of the substance, demonstrated in these works, allowed constructing in [145–154] a model of nanofluid flow in nanochannels, which under certain conditions predicts an unusually high flow rate due to the effect of solid clusters slipping along the nanochannel walls. Along with that, the cases where the channel dimensions were smaller than those of the nanoclusters, the effect of sharp deceleration of such liquids flow was described [146].

Formula	References
$h_p = \alpha \ln \left(\frac{\rho}{T_m - T} + \delta \right)$	[113]
$h_p = \frac{1}{(T_k - T)^{\frac{1}{6}}}$	[4]
$h_p = \frac{1}{(T_m - T)^{\frac{1}{24}}}$	[115]
$h_p = \xi \ln(\gamma T)$	[116]
$h_p = \frac{1}{(T_m - T)^{\frac{1}{3}}}$	[119]
α h_p : $(T_m - T + (T_{m2n} - T)H(T_{m2n} - T)\beta + \delta)^{\frac{1}{3}}$	[11]

TABLE 2. Dependence of the thickness of the non-autonomous phase on temperature

The dependencies of the non-autonomous phase thickness on temperature given in Table 2 describe the equilibrium (locally equilibrium) value of its thickness. However, the dimensional parameters of the particles of a substance in the non-autonomous state may turn out to be larger than those described by these expressions. In such cases, this nonautonomous state can be regarded as metastable, associated with the existence of some spatial or other limitations for mass transfer, described, for instance, in [155]. These limitations on the amount of substance in the system do not allow the formation of an autonomous phase, i.e. a solid phase with a structure that it would have at the given temperature and pressure if the said limitations did not exist. A substance in the metastable non-autonomous state can stably exist in the amorphous or crystalline state, but its crystalline structure will differ from that of the equilibrium (autonomous) phase. Similar cases of metastable non-autonomous states of substance have been experimentally discovered and theoretically described for a large number of oxide compounds and some other substances [94–106]. It should be noted that the thermodynamic estimate of the particle sizes of a substance in some metastable non-autonomous state may differ from the actually observed sizes [89, 156]. This is connected not only with the evaluative nature of the calculated data, but also with the role of kinetic factors, which can play a decisive role at temperatures below the melting point of non-autonomous phases [156].

Another important characteristic of non-autonomous phases for the analysis of interaction processes in solid polycrystalline systems is the viscosity of the liquid (liquid-like) non-autonomous phase. It was shown in [157] that by assuming that the viscosity of a liquid non-autonomous phase at a certain temperature can be calculated by extrapolating the dependency of the viscosity of the molten autonomous phase of the same composition on a given temperature, it is possible to obtain a quantitative match between the experimental and calculated data for the rate of particles migration through ice. Such an approach to predicting the non-autonomous phase viscosity is also corroborated by the temperature dependency of the diffusion coefficients of the substance in it (Fig. 3), since these values are closely related to each other [158]. Based on the above-mentioned nature of the dependency of the non-autonomous phase viscosity on temperature, some phenomena that can be observed during the liquid non-autonomous phase flow in polycrystalline systems were theoretically predicted in [12, 81, 159, 160].

Of considerable interest from the point of view of understanding the structural features of non-autonomous phases are the results of an experimental study of the nature of the viscous behavior of the surface layer (surface – non-autonomous phase) of the crystalline phase, given in [126]. It was experimentally shown in this work that the nature of the viscous behavior of the surface layer at relatively high values of the homologous temperature can only be described if this layer is imagined as inhomogeneous, consisting of a liquid phase with inclusions of solid particles of very small size. It can be noted that essentially this result is close to the conclusions made in [81, 138, 139], in which melting was considered as developing in a rather narrow, but still certain temperature range, as a transition of solid autonomous phase particles to a non-autonomous state, and of the non-autonomous liquid phase to the autonomous one. Another conclusion, apparently very important for understanding the features of the liquid-like non-autonomous phase behavior at high temperatures, is

an experimentally derived conclusion that the thickness of these phases is several times greater compared to the values presented in conclusions of works [161–164].

So to date many basic questions related to the structure and properties of non-autonomous phases remain open. Moreover, this applies both to the region of high values of the homologous temperature, and to the region of temperatures below the melting point of the non-autonomous phase. In the latter case, this is due to the 'frozen' state of the nonautonomous phase, which complicates studies of the processes of approaching this state to equilibrium.

3. Solid-phase reactions and sintering of materials

The experimental works of Tammann and co-authors [107] were aimed at determining the temperature of the onset of active sintering of particles of different composition (T_T) . The further studies, which interpreted the effect discovered in these works as associated with the transition of phase boundaries and grains (non-autonomous phases) into the liquid-like state [119], began considering the Tammann temperature (T_T) as a boundary temperature, upon reaching or exceeding which any processes of solid particles interaction, including solid-phase reactions, start [61]. Along with that, the mechanisms and models of solid-phase chemical reactions, the rate of which was determined by the mass transfer between components, were described as those limited by the rate of diffusion in the solid phase [61, 65]. This situation was apparently due to the fact that the volume fraction of non-autonomous phases in these works was significantly smaller than the fraction of solid (autonomous) phases. The transition of a non-autonomous phase to the liquid-like state could increase only the rate of reagent particles agglomeration. This led to a noticeable increase in the contact area of the interacting particles. In addition, point contacts of solid particles were replaced by areas with a liquid-like interlayer. An increase in the rate of components diffusion through the liquid-like non-autonomous phase in the case of interaction of relatively large particles of micron and even submicron sizes could not significantly affect the total mass transfer rate due to the extremely small proportion of non-autonomous phases in such reaction systems.

The situation changed dramatically when chemical reactions began to take place between nanosized reagent particles. For example, it was shown in [166–169] that in nanosized layers obtained by molecular layering [169–173], no interaction between reagents resulting in formation of solid-phase compounds is observed at temperatures below the non-autonomous phase melting point, and upon reaching this temperature, the synthesis proceeds at such a high rate, which cannot be described as that limited by solid-phase diffusion. A similar situation is observed in the case of synthesizing solid-phase compounds from nanosized reagent powders [174–178]. An analysis of the processes of solid-phase compounds synthesis in oxide compositions consisting of nanosized reagent particles in [179] yields a somewhat paradoxical conclusion that all solid-phase chemical reactions occurring at a high or simply noticeable rate are liquid-phase to some extent, i.e. a liquidlike non-autonomous phase must first form in the reaction system for such reactions to start. The role of non-autonomous phases in the processes of solid particles interaction was considered in more detail in [15]. Numerous examples of the influence of liquid-like non-autonomous phases on mass transfer in solid-phase systems and on the synthesis of solidphase compounds are given in [28–34, 36, 180].

Other processes in solid-phase systems that start only after the transition of the non-autonomous phases to the liquidlike state, as was shown, e.g. in [81, 107, 108], are the processes of sintering and inelastic deformation of ceramics by the mechanism of high-temperature creep. It should be noted that the role of liquid-like non-autonomous phases can consist not only in increasing the rate of particles agglomeration due to the facilitation of their sliding relative to each other due to the appearance of a liquid-like non-autonomous phase, but, as was shown in [9, 10, 12, 81], also by the transformation-transport mechanism of solid-phase particles interaction. Such interactions are schematically illustrated in Fig. 4a,b. It was shown in [81] that proceeding from a more general model relationship and variation of the expressions for the non-autonomous phase viscosity, the transformation-transport mechanism makes it possible to obtain practically the same dependencies as a set of known models of sintering and inelastic deformation of materials obtained by applying different mechanisms. The corresponding comparison of models is given in Table 3.

An important role in solid-phase chemical reactions, in the processes of sintering and inelastic deformation of ceramics, can be played by small additives of relatively fusible solid-phase substances, even in cases where they are chemically inert to other components of the solid-phase system. It was shown in [5,81,182] that such additives get quickly distributed over the surface of other components at a temperature equal to or higher than the melting point of the non-autonomous phase based on such additives. This, in turn, promotes agglomeration of particles of the solid-phase system and, as a consequence, an increase in the rate of solid-phase synthesis and sintering of ceramics.

Another area for which the transition of non-autonomous phases to the liquid-like state is important is the construction of phase equilibrium diagrams in solid-phase systems. For instance, while the approach to the equilibrium state in oxide solid-phase systems at high temperatures requires hours, tens or hundreds of hours [185], the process of approaching equilibrium in the low-temperature region can take a year or more [185], or even geological epochs. In this regard, the analysis of characteristic values of the diffusion rate in solid phases in [185] resulted in proposing the time criterion as a criterion for dividing into high-temperature and low-temperature processes. When the time of approaching the equilibrium state in a solid-phase system is a year and more, then it is proposed to define such systems as low-temperature ones [185]. However, the above analysis of the role of non-autonomous phases in the behavior of solid-phase systems allows proposing another criterion for dividing solid-phase processes into high-temperature and low-temperature ones.

		Average	Mechanical force		
Model	Time, t	particle	per square of the		
		radius	average particle		
		$\hat{R} = \frac{d}{2}$	radius, $f = \frac{F}{\hat{R}}$		
Transformation-transport mechanism (approximation without taking					
into account the contribution of dislocation tubes to transport processes)					
Asymptotic equation of the initial	t^1	\hat{R}^{-1}			
stage of sintering [81]					
Equation of sintering limited by					
mass transfer of a substance through a		$t^{\frac{2}{3}\cdot \frac{n}{1+n}}$ $\hat{R}^{-\frac{2}{3}\cdot \frac{1+2n}{1+n}}$	$f^{\frac{2}{3} \cdot \frac{1}{1+n}}$		
liquid non-autonomous phase [81]					
$n=0.5$	$t^{\frac{1}{2}}$	$\hat{R}^{-\frac{8}{9}}$			
n=1 (Newtonian fluid)	$t^{\frac{1}{3}}$	$\overline{\hat{R}^{-1}}$	$\frac{f^{\frac{4}{9}}}{f^{\frac{1}{3}}}$		
$n=2$	$t^{\frac{4}{9}}$	$\hat{R}^{-\frac{10}{9}}$			
$n=3$	$t^{\frac{1}{2}}$	$\overline{\hat{R}^{-\frac{7}{6}}}$	$f^{\frac{1}{6}}$		
$n \rightarrow \infty$	$t^{\frac{2}{3}}$	$\hat{R}^{-\frac{4}{3}}$	f^0		
Compaction due to redistribution of	t^{1+x}	\hat{R}^0	f^1		
particles in the liquid phase [183]					
Volume diffusion					
$[182]$	$t^{0.49}$	$\hat{R}^{-0.98}$	$f^{0.49}$		
$[183]$	$t^{0.4}$	$\hat{R}^{-0.8}$	$f^{0.40}$		
$[184]$	$t^{\frac{1}{2}}$	\hat{R}^{-1}	$f^{\frac{1}{2}}$		
$[182]$	$t^{0.46}$	$\overline{\hat{R}^{-0.92}}$	$f^{0.46}$		
[116]	$t^{0.48}$	$\hat{R}^{-0.96}$	$f^{0.48}$		
Volume diffusion to grain boundaries	$t^{0.53}$	$\hat{R}^{-1.06}$	$f^{0.53}$		
and spherical surface of particles [184]					
Grain boundaries diffusion					
$[182]$	$t^{\frac{1}{3}}$	\hat{R}^{-1}	$f^{\frac{1}{3}}$		
[184]	$t^{\frac{1}{3}}$	\hat{R}^{-1}	$f^{\frac{1}{3}}$		
[184]	$t^{0.31}$	$\hat{R}^{-0.93}$	$f^{0.32}$		
[116]	$t^{0.32}$	$\hat{R}^{-0.96}$	$f^{0.31}$		
Limited dissolution of the solid phase					
in the liquid and diffusion-viscous	$t^{\frac{1}{3}}$	\hat{R}^{-1}	$f^{\frac{1}{3}}$		
flow of the liquid phase [183]					

TABLE 3. Parametric dependence of shrinkage during sintering and hot pressing of materials

FIG. 4. Schematic representation of the transformation-transport mechanism of solid phase particles interaction

Apparently, it makes more physical sense to use the temperature of the non-autonomous phases transition to the liquidlike state as such a criterion. The corresponding division of processes in solid-phase systems into high-temperature and low-temperature ones is schematically shown in Fig. 5. It should be noted that in terms of the real division of solid-phase systems, the proposed boundary does not differ much from that proposed in [185], but is more justified in physical sense.

FIG. 5. Scheme of processes division in solid-phase systems into high-temperature and low-temperature

4. Conclusion

An analysis of the literature data on the nature of processes in oxide solid-phase systems demonstrated that the decisive role in them is played by the behavior of the substance in the non-autonomous state. A particularly great influence the non-autonomous phases have on the behavior of nanosized systems. The main parameters of the non-autonomous phases that determine the properties and behavior of the solid-phase, especially nanocrystalline ones, are the composition of the non-autonomous phases, the temperature of the non-autonomous phase transition to the liquid-like state, the equilibrium (locally-equilibrium) and metastable thickness of non-autonomous phases, and the viscosity of the liquid-like non-autonomous phase.

The data available in the literature show that there are still a large number of open questions both on the structure and properties of the substance in the non-autonomous state, and on the influence of the non-autonomous state on the behavior and properties of solid-phase systems, which indicates the prospects for further research in this area.

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Submitted 15 September 2024; accepted 28 November 2024

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Conflict of interest: the authors declare no conflict of interest.